



Jacob Schroder

UIUC

The Search for a General Strength of Connection Measure in AMG

CS 550

- 1 Introduction
  - Motivation
  - Smooth Error and Strong Connections
- 2 General Strength of Connection Measure
  - Definition
  - Results
  - Drawbacks
- 3 Conclusion

In an ideal world, AMG should only need a matrix and a smoother, i.e. be a black box solver. That in turn implies that we want a black box strength of connection measure.

The classic Ruge-Stüben, R-S, strength measure covered in class works only for M-matrices. All kinds of problems can generate non M-matrices, such as relatively simple ones like the bi-harmonic or certain anisotropies.

**We need better strength of connection information.**

- Good C-F splitting
- Aggregation done well
- Nonzero placement in the interpolation operator

Algebraically smooth error is the error not reduced by smoothing, the space complementary to the smoother.

$$S = (I - \alpha D^{-1}A).$$

$Se \approx e$  defines smooth error, where

$$\alpha = \frac{K}{\rho(D^{-1}A)} \text{ and } \rho(S) \approx 1.$$

The high energy eigenvectors of  $S$  define the smooth error. This space corresponds to the low energy eigenvectors of  $A$ .

- A strong connection between degrees of freedom  $i$  and  $j$  implies that we can reliably interpolate smooth error between them.
- The chicken and the egg – Defining strong connections assumes some interpolation method, and constructing interpolation operators assumes a definition of smooth error!
- We will concern ourselves only with strength of connection. However, we can say that we want our interpolation basis functions to have low energy, i.e. have a small energy norm (smoothed aggregation). Hence, we may want to define strength using the energy norm.

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The eigenspace of  $S$  defines smooth error, but eigen-calculations are expensive, and it is not clear for an arbitrary operator where to truncate the eigenspace when capturing smooth error.

Can we think of a more accessible heuristic?  $A^{-1}$  relates the size of the residual at dof,  $j$ , to the size of the error at dof,  $i$ .

$$Ae = r \Leftrightarrow A^{-1}r = e. \quad (1)$$

This is the equation MG solves.

Construct a strength of connection matrix,  $\tilde{A}$ , by defining the strength of connection between dof,  $i$ , and dof,  $j$ , as

$$\tilde{A}_{ij} = \frac{\| G_{:,i} - (G_{j,i}) l_{:,j} \|_A}{\| G_{:,i} \|_A}, \quad (2)$$

where  $l_{:,j}$  is the  $j$ -th column of the identity and  $G_{:,i}$  is the  $i$ -th column of  $A^{-1}$  as approximated by some local relaxation scheme such as weighted-Jacobi.

Compare values in  $\tilde{A}_{i,:}$  in order to determine the strong neighbors of dof,  $i$ .

# What Is It Actually Measuring?

- Did the removal of the point from  $G_{:,j}$  create large derivatives?
- If dof,  $i$ , were the only coarse grid point, then a good interpolation operator would be column  $i$  of the inverse (think small energy norm for good basis functions).

If the removal of point  $j$  from  $G_{:,j}$  creates large values in the energy-norm, then points  $i$  and  $j$  are probably strongly connected, because point  $j$ 's removal greatly disrupted the low energy of  $G_{:,j}$ .

Find those points in  $G_{:,j}$  most important to maintaining its *minimal* energy.

# Results - Stencils in $\tilde{A}$ Rotated Anisotropy of $\frac{\pi}{8}$

All results are for the Poisson problem on a regular  $33 \times 33$  grid of bilinear quadrilaterals. Matrices used were non M-matrices.  $\mu$  = the number of weighted-Jacobi steps.

| $\mu$ | Stencil |        |        |
|-------|---------|--------|--------|
| 3     | 0.0006  | 0.0641 | 0.0245 |
|       | 0.0028  | 0.0000 | 0.0028 |
|       | 0.0245  | 0.0641 | 0.0006 |
| 4     | 0.0012  | 0.0852 | 0.0316 |
|       | 0.0018  | 0.0000 | 0.0018 |
|       | 0.0316  | 0.0852 | 0.0012 |

Table:  $\tilde{A}$  stencil at point (17,8)

# Results - Stencils in $\tilde{A}$ Rotated Anisotropy of $\frac{\pi}{4}$

| $\mu$ | Stencil |        |        |
|-------|---------|--------|--------|
| 2     | 0.0011  | 0.0073 | 0.0378 |
|       | 0.0073  | 0.0000 | 0.0073 |
|       | 0.0378  | 0.0073 | 0.0011 |
| 3     | 0.0012  | 0.0151 | 0.0634 |
|       | 0.0151  | 0.0000 | 0.0151 |
|       | 0.0634  | 0.0151 | 0.0012 |
| 4     | 0.0010  | 0.0230 | 0.0840 |
|       | 0.0230  | 0.0000 | 0.0230 |
|       | 0.0840  | 0.0230 | 0.0010 |

Table:  $\tilde{A}$  stencil at point (17,8).

# Results - Stencils in $\tilde{A}$

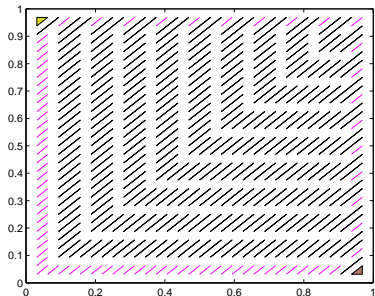
## Coefficient Jump In Center of Domain

| $\mu$ | Point  | Stencil |        |        |
|-------|--------|---------|--------|--------|
| 4     | (17,8) | 0.0356  | 0.0240 | 0.0004 |
|       |        | 0.0469  | 0.0000 | 0.0005 |
|       |        | 0.0356  | 0.0240 | 0.0004 |
| 4     | (18,8) | 0.0003  | 0.0149 | 0.0160 |
|       |        | 0.0005  | 0.0000 | 0.0211 |
|       |        | 0.0003  | 0.0149 | 0.0160 |

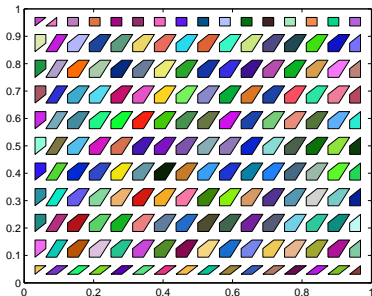
Table:  $\tilde{A}$  stencils around interface.

# Results - Sample Aggregates

## Hot Dog!



Rotated Anisotropy of  $\frac{\pi}{4}$



Rotated Anisotropy of  $\frac{\pi}{8}$

- The core idea is simple and intuitive; look at the entries of  $A^{-1}$  to determine strength of connection.
- Effective for at least some non M-matrices, unlike the classic strength measure.
- Produced MG preconditioners for CG that were as effective as using distance-based strength of connection measures.
- However, the measure is not without its drawbacks.

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What happens when the approximate column of the inverse approaches the actual column of the inverse?

| Stencil |        |        |
|---------|--------|--------|
| 0.1282  | 0.4390 | 0.6246 |
| 0.3854  | 0.0000 | 0.3800 |
| 0.4812  | 0.3162 | 0.0712 |

Table:  $\tilde{A}$  stencil at point (17,8) – Rotated Anisotropy of  $\frac{\pi}{4}$

Perhaps,  $A^{-1}$  with its unrestricted stencil smears information. Use SPAI [4] to generate the approximate inverse,  $G$ . Enforce  $\text{nnz}(G) \approx \text{nnz}(A)$ .

| Stencil |        |        |
|---------|--------|--------|
| 0.0051  | 0.1450 | 0.0042 |
| 0.0095  | 0.0000 | 0.0079 |
| 0.0057  | 0.1599 | 0.0039 |

Table:  $\tilde{A}$  stencil at point (17,8) – Vertically strong anisotropy

SPAI offers limited improvement over direct usage of  $A^{-1}$ .

| Stencil |        |        |
|---------|--------|--------|
| 0.0000  | 0.0374 | 0.0921 |
| 0.0378  | 0.0000 | 0.0285 |
| 0.0944  | 0.0288 | 0.0000 |

Table:  $\tilde{A}$  stencil at point (17,8) –  $\frac{\pi}{4}$  rotated anisotropy

| Stencil |        |        |
|---------|--------|--------|
| 0.0040  | 0.0997 | 0.0227 |
| 0.0005  | 0.0000 | 0.0005 |
| 0.0227  | 0.0997 | 0.0040 |

Table:  $\tilde{A}$  stencil at point (17,8) –  $\frac{\pi}{8}$  rotated anisotropy

## Problem 1 - Cost

$k$  Jacobi relaxations with a 0 initial guess to calculate  $G$  is equivalent to doing  $k - 1$  mat-mats with  $A$ . Calculating the measure increases the cost to  $k$  mat-mats.

## Problem 2 - Not a Panacea

The measure did not perform well for the most difficult test case, a stretched and unstructured grid on an oval-shaped domain. Yet, it did do as well as a distance-based strength of connection measure for this problem.

## Problem 3 – The Real Problem

- $A^{-1}$  relates degrees of freedom with respect to all error for the system,  $Ae = r$ . This includes smooth error, but also error effectively reduced by smoothing.
- Ideally, we want to know for any dof,  $j$ , on the fine grid, what are the most important coarse dof's for interpolating smooth error to  $j$ ?
- This issue is partially avoided by the fact we only relax a handful of times when calculating the approximate inverse. After only a few iterations of relaxation, the approximation to the inverse should still be heavily influenced by the relaxation method and hence, hopefully, local and reflective of the relaxation method.

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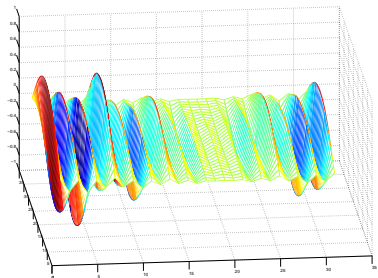
- We need accurate strength of connection information to form coarse grids and interpolation operators in multigrid.
- The proposed measure in [2] is effective for a more general class of matrices than the classic strength of connection measure.
- However, the proposed measure is not without its drawbacks.
  - The cost casts doubt on the computational feasibility of this measure.
  - What happens to the measure in the limit, when using actual columns of  $A^{-1}$ , casts doubt on the idealness of the measure.

- Reduce the cost of the measure by using some further approximation to it or by only calculating for rows that deviate strongly from being “M” like.
- Explore adaptive thresholding by varying the number of weighted-Jacobi iterations.
- Explore other strength measures that are based on approximating the smooth components of  $S$ 's eigenspace.

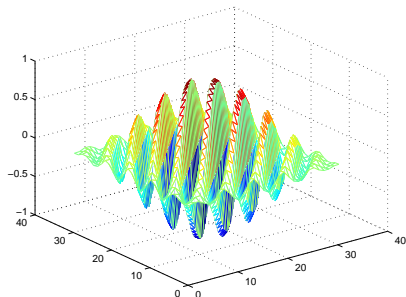
- 1 A. Brandt, S. McCormick and J. W. Ruge. *Algebraic multigrid (AMG) for sparse matrix equations*. Sparsity and Its Applications, DJ Evans, ed., Cambridge University Press, Cambridge, 1984.
- 2 J. Brannick, M. Brezina, S. MacLachlan, T. Manteuffel, S. McCormick, and J. Ruge. *An energy-based AMG coarsening strategy*. Num. Lin. Alg. Appl., 2006: **13**:133-148.
- 3 J. W. Ruge and K. Stüben. *Algebraic Multigrid*, in Multigrid Methods. S. McCormick, ed., SIAM, Philadelphia, 1987, pp. 73-130.
- 4 M. Grote and T. Huckle. *Parallel preconditioning with sparse approximate inverses*. SIAM J. Sci. Comput., 1997: **18**:838-853.

# Dominant Eigenvectors of the Fake Cycle Matrix

Fake Cycle Matrix  $:= (I - P^T P)S$

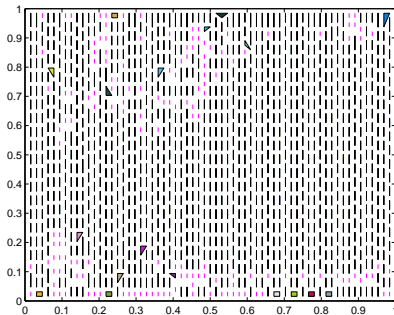


Vertically Strong Anisotropy



Rotated by  $\frac{\pi}{4}$  Anisotropy

# Aggregates For Vertically Strong Anisotropy



Resulting aggregation using eigenvector

$$g_0 = 0 \tag{3}$$

$$g_1 = 0 + \omega D^{-1}(e_j - 0) = \omega D^{-1}e_j$$

$$g_2 = \omega D^{-1}e_j + \omega D^{-1}(e_j - A\omega D^{-1}e_j)$$

Which is  $O(D^{-1}AD^{-1}e_j)$ .

Considering just this dominant operation, the following pattern becomes clear.

$$g_3 = O(D^{-1}AD^{-1}AD^{-1}e_j) = O(A^2e_j) \tag{4}$$

$$g_4 = O(D^{-1}A(D^{-1}AD^{-1}AD^{-1}e_j)) = O(A^3e_j)$$

What if off-diagonals become positive or the matrix loses its definiteness? The classic strength of connection measure begins to break down. For example, the rotated anisotropic diffusion problem,

$$-(\nu c^2 + s^2)u_{xx} + 2(1 - \nu)sc u_{xy} - (\nu s^2 + c^2)u_{yy} = f, \quad (5)$$

with  $\nu = 0.001$ ,  $c = \cos(\frac{\pi}{8})$  and  $s = \sin(\frac{\pi}{8})$  yields a matrix stencil of

$$\begin{array}{ccc} 0.0098 & -0.5200 & -0.3434 \\ 0.1864 & 1.3347 & 0.1864 \\ -0.3434 & -0.5200 & 0.0098 \end{array},$$

for bilinear quads on a structured grid.